

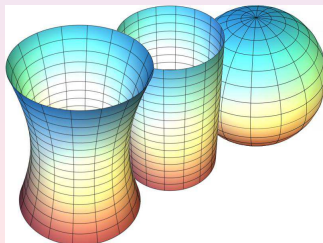
# The Einstein-Hilbert type action on foliated metric-affine manifolds

Vladimir Rovenski

Department of Mathematics, University of Haifa

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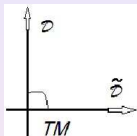
- 1 Pseudo-Riemannian almost-product manifolds
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- $(M^n, g)$  – a **pseudo-Riemannian manifold**
- $\nabla$  – the **Levi-Civita connection**
- **Linear connection**  $\bar{\nabla}$  is  $\nabla$  shifted by **contortion tensor**:  $\mathfrak{T} := \bar{\nabla} - \nabla$ .  
Particular cases of  $\bar{\nabla}$ :
  - (a) Riemann-Cartan geometry: the metric property  $\bar{\nabla} g = 0$ .
  - (b) Torsionless  $\bar{\nabla}$  with symmetric tensor  $\bar{\nabla} g$  arise in statistical structures of differential geometry.
- The **torsion tensor**  $\mathcal{S}$  of  $\bar{\nabla}$  is represented as  $2\mathcal{S}(X, Y) = \mathfrak{T}(X, Y) - \mathfrak{T}(Y, X)$ .
- $\bar{R}_{X,Y} = \bar{\nabla}_Y \bar{\nabla}_X - \bar{\nabla}_X \bar{\nabla}_Y + \bar{\nabla}_{[X,Y]}$  – the **curvature tensor**

# Pseudo-Riemannian almost-product manifolds

Define several tensors for one of distributions; similar tensors for the 2nd distribution are denoted using  $\tilde{\phantom{x}}$  or  $\perp$  notation.

- $\tilde{\mathcal{D}}$  –  $n$ -dimensional **distribution**, **subbundle** of  $TM$
- $\mathcal{D}$  –  $p$ -dimensional distribution,  $g$ -**orthogonal** to  $\tilde{\mathcal{D}}$
- $(\cdot)^\top, (\cdot)^\perp$  – **projections** of  $TM$  onto  $\tilde{\mathcal{D}}$  and  $\mathcal{D}$ .



$\tilde{\mathcal{D}}$ -related tensors,  $X, Y \in \tilde{\mathcal{D}}$ :

- $T(X, Y) = \frac{1}{2} [X, Y]^\perp$  – the **integrability tensor**
- $h(X, Y) = \frac{1}{2} (\nabla_X Y + \nabla_Y X)^\perp$  – **2-nd fundamental form**
- $H = \text{Tr}_g h \in \mathcal{D}$  – the **mean curvature vector**
- $S_{\text{ex}} = g(H, H) - \langle h, h \rangle$  – the **extrinsic scalar curvature**

The  $(1, 1)$ -tensors  $A$  and  $T^\sharp$  correspond to  $h$  and  $T$ :

$$g(A_Z(X), Y) = g(h(X, Y), Z), \quad g(T^\sharp_Z(X), Y) = g(T(X, Y), Z).$$

# The mixed scalar curvature

From now, a distribution  $\tilde{\mathcal{D}} \subset TM$  is fixed.

- $\{E_a \in \Gamma(\tilde{\mathcal{D}}), \mathcal{E}_i \in \Gamma(\mathcal{D})\}$  a local orthonormal frame, and  $\varepsilon_i = g(\mathcal{E}_i, \mathcal{E}_i) = \pm 1$ ,  $\varepsilon_a = g(E_a, E_a) = \pm 1$ .
- $\bar{S}_{\text{mix}} = \frac{1}{2} \sum_{a,i} \varepsilon_a \varepsilon_i [g(\bar{R}(E_a, \mathcal{E}_i)E_a, \mathcal{E}_i) + g(\bar{R}(\mathcal{E}_i, E_a)\mathcal{E}_i, E_a)]$   
– the **mixed scalar curvature**,

$$\bar{S}_{\text{mix}} = S_{\text{mix}} + Q_1 + Q_2, \quad (1)$$

$$S_{\text{mix}} = S_{\text{ex}} + \tilde{S}_{\text{ex}} + \langle T, T \rangle + \langle \tilde{T}, \tilde{T} \rangle + \text{div}(H + \tilde{H}), \quad (2)$$

where (2) is due P. Walczak (1990), and

$$2Q_1 = \sum_{i,a} \varepsilon_i \varepsilon_a [g(\left( (\nabla_i \mathfrak{T})_a - (\nabla_a \mathfrak{T})_i + [\mathfrak{T}_i, \mathfrak{T}_a] \right) E_a, \mathcal{E}_i)],$$

$$2Q_2 = \sum_{i,a} \varepsilon_i \varepsilon_a [g(\left( (\nabla_a \mathfrak{T})_i - (\nabla_i \mathfrak{T})_a + [\mathfrak{T}_a, \mathfrak{T}_i] \right) \mathcal{E}_i, E_a)],$$

If  $\tilde{\mathcal{D}}$  is spanned by a unit vector field  $N$  then  $S_{\text{mix}} = \varepsilon_N \text{Ric}_N$ .

# Flows: the mixed E.-H. action

Let  $\tilde{\mathcal{D}}$  be spanned by a vector field  $N$  with  $\varepsilon_N = g(N, N) = \pm 1$ ,

$\tilde{T}_N^\sharp$  corresponds to the integrability tensor  $\tilde{T}$  of  $\mathcal{D}$ ;

$\tilde{A}_N$  the shape operator of  $\mathcal{D}$ . Set  $\tilde{\tau}_i = \text{Tr } \tilde{A}_N^i$  ( $i = 1, 2$ ).

The **mixed E.-H. action** on  $(M, g, N)$  is defined in [1] by

$$J_{\tilde{\mathcal{D}}, \Omega} : g \mapsto \int_{\Omega} \left\{ \frac{1}{2\alpha} (\varepsilon_N \text{Ric}_N - 2\Lambda) + \mathcal{L}(g) \right\} d \text{vol}_g . \quad (3)$$

$\Lambda$  – a cosmological constant,  $\alpha$  – a coupling constant,  $\mathcal{L} = \mathcal{L}(g)$  – Lagrangian describing the matter contents,

$\Omega \subset M$  – a relatively compact domain in  $(M, g)$ .

**Example (A Lorentzian manifold  $(M, g)$ ,  $\dim M = 4$ ):**

(i)  $(M^4, g)$  with a timelike vector field  $N$  is called a **space-time**.

(ii) **Globally hyperbolic space-times**  $(M^4, g)$  come equipped with a codimension-one foliation, since  $\exists$  time function  $\mathbb{T} : M \rightarrow \mathbb{R}$ ,  $M = \mathbb{R} \times \hat{S}$ , where for each  $t$ ,  $\{t\} \times \hat{S}$  is a Cauchy's surface.

Let  $\tilde{T} = 0$ . The **mixed Ricci tensor** is defined in [1] by

$$\text{Ric}_{\mathcal{D}} := \text{Ric}_{\mathcal{D}}^{\perp} + \text{Ric}'_{\mathcal{D}} + \text{Ric}_{\mathcal{D}}^{\top},$$

where

$$(\text{Ric}_{\mathcal{D}})^{\perp}(X, Y) = (\nabla_N \tilde{h}_{sc} - \tilde{\tau}_1 \tilde{h}_{sc})(X, Y), \quad X, Y \perp N,$$

$$(\text{Ric}_{\mathcal{D}})'(X, N) = \text{Ric}_{\mathcal{D}}^1(g)(N, X) = \text{div}(\tilde{A}_N(X)),$$

$$(\text{Ric}_{\mathcal{D}})^{\top}(N, N) = -\text{div} H,$$

and  $\tilde{h}_{sc} = \varepsilon_N \langle \tilde{h}, N \rangle$ . Put  $\text{Scal}_{\mathcal{D}} := \text{Tr}_g(\text{Ric}_{\mathcal{D}})$ .

**Theorem** ( $\varepsilon_N = -1$ , see [1])

Let  $\tilde{\mathcal{D}} = \text{Span}(N)$  with integrable  $\mathcal{D}$ . If  $g$  is critical for (3) then

$$\text{Ric}_{\mathcal{D}} - \frac{1}{2} \text{Scal}_{\mathcal{D}} \cdot g - \text{Ric}_N (N^b \otimes N^b + \frac{1}{2} g) + \Lambda g = \mathfrak{a} \Theta. \quad (4)$$

Here  $\Theta$  – **stress-energy tensor**,  $\Theta_{ij} = -2 \frac{\partial \mathcal{L}}{\partial g_{ij}} + \mathcal{L} g_{ij}$ .

(A **Riemannian flow** of  $N$  obeys (4) with  $\Theta = \Lambda = 0$ ).

## Goals of:

[1] E. Barletta, S. Dragomir, V. Rovenski and M. Soret:

**Mixed gravitational field equations on globally hyperbolic spacetimes**, Class. Quantum Grav., 30 (2013), 26 pp.

- sufficient conditions, e.g. isoparametric time-function and totally umbilical  $\mathcal{D}$ , for the existence of solutions to (4) in empty space.
- the conservation law, analog of  $\operatorname{div} \Theta = 0$  in relativity.
- linearization and solution of (4) for empty space.
- equations of motion (geodesics) of a particle in the ‘gravitational field’ governed by (4).
- after the linearization of motion equations about the Minkowski metric, the coupling constant  $\alpha = 2\pi k/c^2$  has been determined, where  $k$  is the gravitational constant.



# The mixed Einstein-Hilbert action

$M$  with a smooth distribution  $\tilde{\mathcal{D}}$  of arbitrary dimension.

The **mixed Einstein-Hilbert action** on  $(M, g, \tilde{\mathcal{D}}, \bar{\nabla} = \nabla + \mathfrak{T})$

$$J_{\tilde{\mathcal{D}}, \Omega} : (g, \mathfrak{T}) \mapsto \int_{\Omega} \left\{ \frac{1}{2\alpha} (S_{\text{mix}} - 2\Lambda) + \mathcal{L}(g, \mathfrak{T}) \right\} d \text{vol}_g. \quad (5)$$

$J_{\text{mix}}(g) = \int_{\Omega} S_{\text{mix}} d \text{vol}_g$  – the **gravitational part** of the action.

Question:

- Euler-Lagrange equations of  $J_{\tilde{\mathcal{D}}, \Omega}$ ?
- their solutions: critical metrics and connections ?

A **variation of connection** has the form  $\bar{\nabla}^t = \nabla + \mathfrak{T}^t$ ,  $|t| < \varepsilon$ .

## Proposition (E.-L. equations)

Let  $(M, \tilde{\mathcal{D}}, \mathcal{D}, g)$  be a pseudo-Riemannian almost product manifold. Then the E.-L. equations for (5) with  $\Lambda = \Theta = 0$ , as a functional on the space of all connections on  $M$ , are the following:

$$-2\tilde{T}(\mathcal{E}_i, \mathcal{E}_j) - (\mathfrak{T}_j^* \mathcal{E}_i + \mathfrak{T}_i \mathcal{E}_j)^\top = 0,$$

$$\sum_i \varepsilon_i (\mathfrak{T}_i^* \mathcal{E}_i)^\top = \tilde{H}, \quad \text{for } n > 1,$$

$$\sum_i \varepsilon_i (\mathfrak{T}_i \mathcal{E}_i)^\top = -\tilde{H}, \quad \text{for } n > 1,$$

$$\mathfrak{T}_i^\top - (\mathfrak{T}_i^\top)^* = 2T_i^\sharp,$$

$$\sum_j \varepsilon_j g((\mathfrak{T}_j + \mathfrak{T}_j^*) \mathcal{E}_j, \mathcal{E}_i) \text{Id}_{\tilde{\mathcal{D}}} - \mathfrak{T}_i^\top - (\mathfrak{T}_i^\top)^* = 0,$$

$$h = (1/n) H g^\top,$$

$$\sum_j \varepsilon_j (\mathfrak{T}_j \mathcal{E}_j - \mathfrak{T}_j^* \mathcal{E}_j)^\perp = \frac{2}{n} (n-1) H,$$

and 7 dual equations.

## Proposition

The Levi-Civita connection of a metric  $g \in \text{Riem}(M, \tilde{\mathcal{D}}, \mathcal{D})$  is critical for (5) as a functional on the space of all linear connections on  $M$  if and only if

$$g(\tilde{H}, E_b \delta_{ac} - E_c \delta_{ab}) = \alpha S_{bc;a}, \quad (6a)$$

$$g(H, \mathcal{E}_j \delta_{ik} - \mathcal{E}_k \delta_{ij}) = \alpha S_{jk;i}, \quad (6b)$$

$$-g(h(E_a, E_b) - H \delta_{ab}, \mathcal{E}_i) = \alpha (S_{bi;a} + S_{ab;i}/2), \quad (6c)$$

$$-g(\tilde{h}(\mathcal{E}_i, \mathcal{E}_j) - \tilde{H} \delta_{ij}, E_a) = \alpha (S_{ja;i} - S_{ij;a}/2), \quad (6d)$$

$$2g(\tilde{T}_a^\sharp \mathcal{E}_j, \mathcal{E}_i) = \alpha S_{ij;a}, \quad (6e)$$

$$2g(T_i^\sharp E_b, E_a) = \alpha S_{ab;i}, \quad (6f)$$

and the spin tensor  $S$  should satisfy the following compatibility conditions when  $\alpha \neq 0$ :

$$S_{bc;a} = -S_{cb;a}, \quad S_{jk;i} = -S_{kj;i}, \quad S_{ja;i} = -S_{aj;i}, \quad S_{bi;a} = -S_{ib;a}, \quad (7a)$$

$$\sum_a (S_{ai;a} + S_{aa;i}/2) = 0, \quad \sum_i (S_{ia;i} + S_{ii;a}/2) = 0. \quad (7b)$$

## Example

Consider a spin fluid  $S_{\mu\nu}^\eta(x) = \lambda^\eta(x) s_{\mu\nu}$ , with skew-symmetric spin:  $s_{\mu\nu} = -s_{\nu\mu}$ . Obviously, (7a) are satisfied, and (7b) reads

$$\sum_a \lambda_a s_{ai} = 0, \quad \sum_i \lambda_i s_{ai} = 0,$$

from which nonzero  $\lambda = (\lambda_\mu)$  can be found when  $\dim M > 2$ .

One may examine smaller spaces of varied connections, e.g., metric connections, corresponding to Riemann-Cartan manifolds.

## Corollary

*Let  $n + p > 2$ . The Levi-Civita connection of  $g$  is critical for the gravitational part of (5) on the space of all metric connections if and only if both distributions are totally geodesic and integrable, i.e.,  $M$  splits.*

# Variations of metric

From now assume  $\mathfrak{T} = 0$ .

We consider variations  $\{g_t \in \text{Riem}(M) : |t| < \varepsilon\}$  of metric  $g_0 = g$  on  $M$  such that the induced infinitesimal variations, presented by a symmetric  $(0, 2)$ -tensor  $B_t \equiv \partial g_t / \partial t$ , are supported in a relatively compact domain  $\Omega$  in  $M$ .

If  $\tilde{\mathcal{D}}$  and  $\mathcal{D}$  are fixed and their orthogonality is preserved by  $(g_t)$  then we have **adapted variation** of metric, see [2, 3]. Adapted variations preserving metric on  $\mathcal{D}$  are called  **$g^\top$ -variations**.

A tensor  $B = \partial_t g_t \in \text{Sym}^2(M)$  is orthogonally decomposed

$B = B^\perp + B' + B^\top$ , with symmetric tensors

$$B^\top(X, Y) := B(X^\top, Y^\top),$$

$$B^\perp(X, Y) := B(X^\perp, Y^\perp),$$

$$B'(X, Y) := \frac{1}{2}(B(X^\top, Y^\perp) + B(X^\perp, Y^\top)).$$

## Lemma

Let a local  $(\tilde{\mathcal{D}}, \mathcal{D})$ -adapted and orthonormal for  $t = 0$  frame  $\{E_a(t), \mathcal{E}_i(t)\}$  be evolved by  $g_t \in \text{Riem}(M)$ ,  $|t| < \varepsilon$ , according to

$$\partial_t E_a = -(1/2) B_t^\sharp(E_a)^\top, \quad \partial_t \mathcal{E}_i = -\frac{1}{2} B_t^\sharp(\mathcal{E}_i)^\perp - B_t^\sharp(\mathcal{E}_i)^\top. \quad (8)$$

Then, for all  $t$ ,  $\{E_a(t), \mathcal{E}_i(t)\}$  is a  $g_t$ -orthonormal frame adapted to  $\tilde{\mathcal{D}}$ , i.e.,  $\{E_a(t)\}$  belong to  $\tilde{\mathcal{D}}$ .

## Definition

Define symmetric  $(0, 2)$ -tensors  $\alpha, \theta, \tilde{\delta}$  by

$$2\alpha(X, Y) = A_{X^\perp}(Y^\top) + A_{Y^\perp}(X^\top),$$

$$2\theta(X, Y) = T_{X^\perp}^\#(Y^\top) + T_{Y^\perp}^\#(X^\top),$$

$$2\tilde{\delta}_Z(X, Y) = g(\nabla_{X^\top} Z, Y^\perp) + g(\nabla_{Y^\top} Z, X^\perp).$$

Define  $(1, 1)$ -tensors  $\mathcal{T} := \sum_i \varepsilon_i (T_i^\#)^2$ ,  $\mathcal{K} = \sum_i \varepsilon_i [T_i^\#, A_i]$ ,  
and a  $(0, 2)$ -tensor  $\operatorname{div} h = \sum_\nu \varepsilon_\nu g(\nabla_\nu h, e_\nu)$ .

For any  $(0, 2)$ -tensors  $P, Q$  and  $S$ , define a tensor  $\Upsilon_{P, Q}$  by

$$\langle \Upsilon_{P, Q}, S \rangle = \sum_{\lambda, \nu} \varepsilon_\lambda \varepsilon_\nu [S(P(e_\lambda, e_\nu), Q(e_\lambda, e_\nu)) + S(Q(e_\lambda, e_\nu), P(e_\lambda, e_\nu))]$$

where  $\{e_\lambda\}$  is an orthonormal basis and  $\varepsilon_\lambda = g(e_\lambda, e_\lambda) \in \{-1, 1\}$ .

Proposition ( For any variation  $g_t$  we have )

$$\partial_t \langle \tilde{h}, \tilde{h} \rangle = \langle \operatorname{div} \tilde{h} - 4\Upsilon_{\tilde{\alpha}, \theta} + \tilde{\mathcal{K}}^b - (1/2) \Upsilon_{\tilde{h}, \tilde{h}}, B \rangle - \operatorname{div} \langle \tilde{h}, B \rangle,$$

$$\partial_t g(\tilde{H}, \tilde{H}) = \langle (\operatorname{div} \tilde{H}) g^\perp + 4\langle \theta, \tilde{H} \rangle - \tilde{H}^b \otimes \tilde{H}^b, B \rangle - \operatorname{div}((\operatorname{Tr}_{\mathcal{D}} B^\sharp) \tilde{H}),$$

$$\begin{aligned} \partial_t \langle h, h \rangle = & \langle \operatorname{div} h + \mathcal{K}^b - 2(\operatorname{div} \alpha)_{|\tilde{\mathcal{D}} \times \mathcal{D}} - 2(\operatorname{div} \alpha)_{|\mathcal{D} \times \tilde{\mathcal{D}}} - 2\Upsilon_{\alpha, \tilde{\alpha} + \tilde{\theta}} \\ & - (1/2) \Upsilon_{h, h}, B \rangle + 2 \operatorname{div} \langle \alpha, B \rangle - \operatorname{div} \langle h, B \rangle, \end{aligned}$$

$$\begin{aligned} \partial_t g(H, H) = & \langle 2\langle \tilde{\theta} - \tilde{\alpha}, H \rangle + 2 \operatorname{Sym}(H^b \otimes \tilde{H}^b) - 2\tilde{\delta}_H + (\operatorname{div} H) g^\top \\ & - H^b \otimes H^b, B \rangle + 2 \operatorname{div}((B^\sharp H)^\top) - \operatorname{div}((\operatorname{Tr}_{\tilde{\mathcal{D}}} B^\sharp) H), \end{aligned}$$

$$\begin{aligned} \partial_t \langle \tilde{T}, \tilde{T} \rangle = & \langle 2\tilde{\mathcal{T}}^b + 2\Upsilon_{\tilde{\theta}, \theta - \alpha} - 2(\operatorname{div} \tilde{\theta})_{|\tilde{\mathcal{D}} \times \mathcal{D}} - 2(\operatorname{div} \tilde{\theta})_{|\mathcal{D} \times \tilde{\mathcal{D}}} \\ & + (1/2) \Upsilon_{\tilde{T}, \tilde{T}}, B \rangle + 2 \operatorname{div} \langle \tilde{\theta}, B \rangle, \end{aligned}$$

$$\partial_t \langle T, T \rangle = \langle (1/2) \Upsilon_{T, T} + 2\mathcal{T}^b, B \rangle.$$

The formulas are applicable for various curvature functionals on a pseudo-Riemannian manifold with a distribution, e.g. for the **total extrinsic scalar curvature** of  $\tilde{\mathcal{D}}$ .



## Definition

The symmetric  $(0, 2)$ -tensor  $\text{Ric}_{\mathcal{D}} = (\text{Ric}_{\mathcal{D}})^{\perp} + (\text{Ric}_{\mathcal{D}})^{\prime} + (\text{Ric}_{\mathcal{D}})^{\top}$  with the trace  $\text{Scal}_{\mathcal{D}} = \text{Tr}_g(\text{Ric}_{\mathcal{D}})$  and the components

$$(\text{Ric}_{\mathcal{D}})^{\perp} = \text{div } \tilde{h} + \tilde{\mathcal{K}}^b - 2\tilde{\mathcal{T}}^b + H^b \otimes H^b - \frac{1}{2}\Upsilon_{h,h} - \frac{1}{2}\Upsilon_{T,T} + \eta g^{\perp},$$

$$\begin{aligned} \frac{1}{2}(\text{Ric}_{\mathcal{D}})^{\prime} &= \tilde{\delta}_H + (\text{div}(\tilde{\theta} - \alpha))|_{\tilde{\mathcal{D}} \times \mathcal{D}} + (\text{div}(\tilde{\theta} - \alpha))|_{\mathcal{D} \times \tilde{\mathcal{D}}} - 2\langle \theta, \tilde{H} \rangle \\ &\quad - \langle \tilde{\theta} - \tilde{\alpha}, H \rangle - 2\Upsilon_{\tilde{\alpha}, \theta} - \Upsilon_{\alpha, \tilde{\alpha}} - \Upsilon_{\tilde{\theta}, \theta} - \text{Sym}(H^b \otimes \tilde{H}^b), \end{aligned}$$

$$(\text{Ric}_{\mathcal{D}})^{\top} = \text{div } h + \mathcal{K}^b - 2\mathcal{T}^b + \tilde{H}^b \otimes \tilde{H}^b - \frac{1}{2}\Upsilon_{\tilde{h}, \tilde{h}} - \frac{1}{2}\Upsilon_{\tilde{T}, \tilde{T}} + \mu g^{\top}$$

is called the **mixed Ricci curvature** of  $g \in \text{Riem}(M, \tilde{\mathcal{D}}, \mathcal{D})$ . Here

$$\eta = -\frac{n-1}{p+n-2} \text{div}(\tilde{H} - H), \quad \mu = \frac{p-1}{p+n-2} \text{div}(\tilde{H} - H). \quad (9)$$

## Theorem (Mixed gravitational field equations)

A metric  $g \in \text{Riem}(M, \tilde{\mathcal{D}}, \mathcal{D})$  is critical for the action (5) if and only if  $g$  is a solution to the Euler-Lagrange equations

$$\text{Ric}_{\mathcal{D}} - (1/2) \text{Scal}_{\mathcal{D}} \cdot g + \Lambda g = \alpha \Theta. \quad (10)$$

The following symmetric  $(0, 2)$ -tensor:

$$G_{\mathcal{D}} := \text{Ric}_{\mathcal{D}} - (1/2) \text{Scal}_{\mathcal{D}} \cdot g$$

is referred to as the **mixed Einstein tensor**. Note that

$$\text{Tr}_g G_{\mathcal{D}} = \frac{2 - n - p}{2} \text{Scal}_{\mathcal{D}},$$

and for  $n = p = 1$  we have  $G_{\mathcal{D}} = 0$ .

## Corollary (The mixed field equations for a flow)

Let  $\tilde{\mathcal{D}}$  be spanned by a unit vector field  $N$ . Then  $g$  is critical for the action (3) if and only if (10) holds, or, in components,

$$\begin{aligned} \nabla_N \tilde{h}_{sc} - \tilde{\tau}_1 \tilde{h}_{sc} - \varepsilon_N (2(\tilde{T}_N^\sharp)^2 + [\tilde{A}_N, \tilde{T}_N^\sharp])^\flat + \langle \tilde{T}, \tilde{T} \rangle \\ - \frac{1}{2} (\varepsilon_N (2N(\tilde{\tau}_1) - \tilde{\tau}_1^2 - \tilde{\tau}_2) - 2\Lambda) g^\perp = \mathfrak{a} \Theta|_{\mathcal{D} \times \mathcal{D}}, \end{aligned}$$

$$(\operatorname{div}^\perp \tilde{T}_N^\sharp)^\sharp + 2 \tilde{T}_N^\sharp(H) = \mathfrak{a} \Theta|_{\mathcal{D} \times \tilde{\mathcal{D}}},$$

$$(1/2) \varepsilon_N (\tilde{\tau}_1^2 - \tilde{\tau}_2) - (3/2) \langle \tilde{T}, \tilde{T} \rangle + \Lambda = \mathfrak{a} \Theta_{N,N},$$

where  $\Theta_{N,N} : M \rightarrow \mathbb{R}$  is defined by  $\Theta|_{\tilde{\mathcal{D}} \times \tilde{\mathcal{D}}} = \Theta_{N,N} g^\top$ .

These equations for a **space-time** generalize result in [1] for a **globally hyperbolic space-time** (i.e.,  $\tilde{T} = 0$ ).

# Flows: with integrable normal distribution

For the flow of  $N$  with integrable normal distribution, the mixed Ricci tensor  $\text{Ric}_{\mathcal{D}}(g)$  has components

$$\begin{aligned}(\text{Ric}_{\mathcal{D}})^{\perp}(g) &= \nabla_N \tilde{h}_{sc} - \tilde{\tau}_1 \tilde{h}_{sc}, \\(\text{Ric}_{\mathcal{D}})^{\prime}(g) &= 0, \\(\text{Ric}_{\mathcal{D}})^{\top}(g) &= \varepsilon_N(N(\tilde{\tau}_1) - \tilde{\tau}_2) g^{\top}.\end{aligned}\tag{11}$$

Comparing (10):

$$\text{Ric}_{\mathcal{D}} - \frac{1}{2} \text{Scal}_{\mathcal{D}} \cdot g + \Lambda g = \alpha \Theta$$

with (4) of [1]:

$$\text{Ric}_{\mathcal{D}} - \frac{1}{2} \text{Scal}_{\mathcal{D}} \cdot g - \text{Ric}_N(N^b \otimes N^b + \frac{1}{2} g) + \Lambda g = \alpha \Theta,$$

we see that (9) is the best choice of  $(\eta, \mu)$  in definition of the mixed Ricci tensor.

# Biconformal solutions of the E.-L. equations

Proposition (Double-twisted products,  $\phi, \psi \in C^\infty(M_1 \times M_2)$ )

Let  $\bar{g} = \psi g^\top + \phi g^\perp$ , where  $g = g^\top + g^\perp$  is a product metric.  
Then  $\bar{g}$  is critical for the action (5) if and only if the following hold:

$$(\mathbf{F}_1 + \Lambda) g^\perp + \mathbf{F}_2 (\nabla^\perp \phi)^b \otimes (\nabla^\perp \phi)^b = \mathbf{a} \Theta|_{\mathcal{D} \times \mathcal{D}},$$

$$\mathbf{F}_3 (\nabla^\top \phi)^b \otimes (\nabla^\perp \psi)^b + \mathbf{F}_4 (\text{Hess}_\psi)|_{\tilde{\mathcal{D}} \times \mathcal{D}} = \mathbf{a} \Theta|_{\tilde{\mathcal{D}} \times \mathcal{D}},$$





$$(\mathbf{F}_5 + \Lambda) g^\top + \mathbf{F}_6 (\nabla^\top \phi)^b \otimes (\nabla^\top \phi)^b = \mathbf{a} \Theta|_{\tilde{\mathcal{D}} \times \tilde{\mathcal{D}}},$$

where  $F_5$  and  $F_6$  are 'dual' to  $F_1$  and  $F_2$ , and

$$\mathbf{F}_1 = \frac{n(n-1)}{4} \psi^{-2} (\nabla^\perp \psi)^2 + \frac{4-p}{4} (p-1) \psi^{-1} \phi^{-1} (\nabla^\top \phi)^2 \\ - (p-1) \psi^{-1} \Delta^\top \phi + \frac{2-n}{2} (p-1) \psi^{-2} g(\nabla^\top \psi, \nabla^\top \phi),$$

$$\mathbf{F}_3 = -\frac{n-1}{2} \psi^{-1} \left( \frac{p-2}{2} \phi^{-1} - \psi^{-1} \right),$$

$$\mathbf{F}_2 = -\frac{n(n-1)}{2} \phi^{-2} \psi^{-2}, \quad \mathbf{F}_4 = \frac{1-n}{2} \psi^{-1}.$$

-  [1] E. Barletta, S. Dragomir, V. Rovenski and M. Soret: **Mixed gravitational field equations on globally hyperbolic spacetimes**, Class. Quantum Grav., 30 (2013), 26 pp.
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# THANK YOU !

